



B. L. D. E. Association's

S. B. ARTS & K. C. P. SCIENCE COLLEGE, VIJAYAPUR

1st

INTERNAL TEST 2023 -2024

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20 / 20

3a] Schrodinger time independent wave equation

$$\text{W.K.T } \psi(x,t) = A e^{-\frac{2\pi i}{h}(Et - Px)}$$

$$\psi(x,t) = A e^{-\frac{2\pi i}{h}(Et - Px)} \quad \text{--- (1)}$$

$$\psi(x,t) = A e^{-\frac{2\pi i E}{h}t} \cdot e^{-\frac{2\pi i P}{h}x} \quad \text{--- (2)}$$

$$\psi(x,t) = \psi_0 e^{-\frac{2\pi i E}{h}t}$$

$$\therefore \psi_0 = A e^{-\frac{2\pi i P}{h}x}$$

$$\psi(x,t) = \psi_0 e^{-\frac{2\pi i E}{h}t} \quad \text{--- (3)}$$

Eqn (3) diff w.r to 't'

$$\frac{\partial \psi(x,t)}{\partial t} = \psi_0 e^{-\frac{2\pi i E}{h}t} \cdot \left(-\frac{2\pi i E}{h}\right)$$

$$\frac{\partial \psi}{\partial t} = \psi_0 \left(-\frac{2\pi i E}{h}\right) \quad \therefore \psi_0 e^{-\frac{2\pi i E}{h}t} = \psi$$

$$\frac{\partial \psi}{\partial t} = -\psi \frac{2\pi i E}{h} \quad \text{--- (4)}$$

Eqn (3) diff w.r to 'x'

$$\frac{\partial \psi}{\partial x} = \frac{\partial \psi_0}{\partial x} e^{-\frac{2\pi i E}{h}t}$$

$$\frac{\partial^2 \psi}{\partial x^2} = \frac{\partial^2 \psi_0}{\partial x^2} e^{-\frac{2\pi i E}{h}t} \quad \text{--- (5)}$$

consider Schrodinger time independent eqn

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V\psi$$

Eqn (4) & (5)

$$i\hbar \left[-\psi \frac{2\pi i E}{h}\right] = \left[-\frac{\hbar^2}{2m} \left(\frac{\partial^2 \psi_0}{\partial x^2} e^{-\frac{2\pi i E}{h}t}\right)\right] + V\psi$$

$$E\psi = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V\psi$$

$$\therefore \psi_0 e^{-\frac{2\pi i E}{h}t} = \psi$$

In 1-D

$$E\psi = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V\psi$$

In 3-D

$$E\psi = -\frac{\hbar^2}{2m} \nabla^2 \psi + V\psi$$

3b) Statement :- consider two Eigen functions ψ_m & ψ_n
 these two f's are said to be orthonormalized
 i.e. $\int_{-\infty}^{\infty} \psi_m^* \psi_n = \delta_{mn}$ where $\delta_{mn} = \begin{cases} 0 & m \neq n \\ 1 & m = n \end{cases}$

proof :- consider Schrodinger time independent Eqⁿ & is given by

$$E\psi = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V\psi = \text{--- (1)}$$

Eqⁿ (1) can be written as ψ_m & ψ_n of the wave fⁿ

$$E_m \psi_m = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi_m}{\partial x^2} + V\psi_m$$

$$E_m \psi_m = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi_m}{\partial x^2} + V\psi_m \text{ --- (2)}$$

$$E_n \psi_n = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi_n}{\partial x^2} + V\psi_n \text{ --- (3)} \quad \times \psi_m^*$$

It's complex conjugate the above Eqⁿs.

$$E_m \psi_m^* = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi_m^*}{\partial x^2} + V\psi_m^* \text{ --- (4)} \quad \times \psi_n$$

$$E_n \psi_n^* = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi_n^*}{\partial x^2} + V\psi_n^* \text{ --- (5)}$$

from Eqⁿ (2) multiply by ψ_m^* & Eqⁿ (3) multiply by ψ_n
 by left then subtract

$$\psi_n E_m \psi_m^* - \psi_m^* E_n \psi_n = \left[-\frac{\hbar^2}{2m} \psi_n \frac{\partial^2 \psi_m^*}{\partial x^2} + \psi_n V \psi_m^* \right] - \left[-\frac{\hbar^2}{2m} \psi_m^* \frac{\partial^2 \psi_n}{\partial x^2} + \psi_m^* V \psi_n \right]$$

$$(E_m - E_n) \psi_n \psi_m^* = -\frac{\hbar^2}{2m} \left[\psi_n \frac{\partial^2 \psi_m^*}{\partial x^2} - \psi_m^* \frac{\partial^2 \psi_n}{\partial x^2} \right]$$

$$(E_m - E_n) \psi_n \psi_m^* = -\frac{\hbar^2}{2m} \left[\psi_n \frac{\partial^2 \psi_m^*}{\partial x^2} - \psi_m^* \frac{\partial^2 \psi_n}{\partial x^2} \right]$$

$$\frac{2m}{\hbar^2} [E_m - E_n] \psi_n \psi_m^* = - \left[\frac{\psi_n \partial^2 \psi_m^*}{\partial x^2} - \frac{\psi_m^* \partial^2 \psi_n}{\partial x^2} \right]$$

$$\frac{2m}{\hbar^2} [E_m - E_n] \psi_n \psi_m^* = \left[\frac{\psi_m^* \partial^2 \psi_n}{\partial x^2} - \frac{\psi_n \partial^2 \psi_m^*}{\partial x^2} \right]$$

Integration on b.s betⁿ the limit $+\infty$ to $-\infty$

$$\frac{2m}{\hbar^2} (E_m - E_n) \int_{-\infty}^{\infty} \psi_n \psi_m^* dx = \int_{-\infty}^{\infty} \left(\psi_m^* \frac{\partial^2 \psi_n}{\partial x^2} - \psi_n \frac{\partial^2 \psi_m^*}{\partial x^2} \right) dx$$

$$\frac{2m}{\hbar^2} (E_m - E_n) \int_{-\infty}^{\infty} \psi_n \psi_m^* dx = \left[\psi_m^* \frac{\partial \psi_n}{\partial x} - \psi_n \frac{\partial \psi_m^*}{\partial x} \right]_{-\infty}^{\infty}$$

$\psi_m^* \frac{\partial \psi_n}{\partial x}, \frac{\partial \psi_m^*}{\partial x} \text{ vanishes}$

$$\frac{2m}{\hbar^2} (E_m - E_n) \int_{-\infty}^{\infty} \psi_n \psi_m^* dx = 0$$

$$\frac{2m}{\hbar^2} (E_m - E_n) \neq 0$$

$$\int_{-\infty}^{\infty} \psi_m^* \psi_n dx = 0$$

$$m \neq n$$

$$i) \int_{-\infty}^{\infty} \psi_m^* \psi_n dx = 1$$

$$m \neq n$$

$$ii) \int_{-\infty}^{\infty} \psi_m^* \psi_n dx = 0$$

a) Schrodinger time dependent Eqn $\psi(x,t) = A e^{i(kx - \omega t)}$
 Let $\psi(x,t) = A e^{-2\pi i (Et - x/\lambda)}$ — (1)
 $E = h\nu$, $\lambda = h/p$

$$\psi(x,t) = A e^{-\frac{2\pi i}{h} (Et - Px)} \quad \text{--- (2)}$$

Eqn (2) diff w.r to 't'.

$$\frac{\partial \psi}{\partial t} = A e^{-\frac{2\pi i}{h} Et} \left(-\frac{2\pi i E}{h} \right)$$

$$\frac{\partial \psi}{\partial t} = \psi \left(-\frac{2\pi i E}{h} \right) = -\frac{2\pi i E}{h} \psi \quad \text{--- (3)}$$

Eqn (2) diff w.r to 'x'

$$\frac{\partial \psi}{\partial x} = A e^{-\frac{2\pi i}{h} Px} \cdot \left(-\frac{2\pi i P}{h} \right)$$

$$\frac{\partial^2 \psi}{\partial x^2} = A e^{-\frac{2\pi i}{h} Px} \left(-\frac{2\pi i P}{h} \right)^2$$

$$\frac{\partial^2 \psi}{\partial x^2} = \psi \left(-\frac{2^2 \pi^2 P^2}{h^2} \right)$$

$$\frac{\partial^2 \psi}{\partial x^2} = -\frac{4\pi^2 P^2}{h^2} \psi \quad \text{--- (4)}$$

Total Energy is equal to the sum of K.E & Potential Energy

$$E = \frac{p^2}{2m} + V(x,t) \quad \text{--- (5)}$$

Eq (5) multiply by ψ

$$E\psi = \frac{p^2}{2m}\psi + V\psi \quad \text{--- (7)}$$

from Eq (3) & (6) substitute in Eq (7)

$$\frac{-h}{2\pi i} \frac{\partial \psi}{\partial t} = \frac{h^2}{2m} \cdot \left(-\frac{1}{4\pi^2} \frac{\partial^2 \psi}{\partial x^2} \right) + V\psi$$

$$\frac{-h}{2\pi i} \frac{\partial \psi}{\partial t} = \frac{h^2}{8\pi^2 m} \frac{\partial^2 \psi}{\partial x^2} + V\psi$$

$$\boxed{i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V\psi}$$

It is S. time dependent 1-D Eq.

$$\boxed{i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + V\psi}$$

It is 3-D Eq. S.D.E.

b) i) $[z, P_z]$ $\Rightarrow [z, -i\hbar \frac{\partial}{\partial z}] = -i\hbar [z, \frac{\partial}{\partial z}]$
 $z = z \quad P_z = -i\hbar \frac{\partial}{\partial z}$

consider $[z, P_z] f(x)$

$$[z, P_z] f(x) = [z P_z - P_z z] f(x)$$

$$= [z (-i\hbar \frac{\partial}{\partial z}) - (-i\hbar \frac{\partial}{\partial z}) z] f(x)$$

$$= [-z i\hbar \frac{\partial}{\partial z} + i\hbar \frac{\partial z}{\partial z}] f$$

$$= [-z i\hbar \frac{\partial f}{\partial z} + i\hbar \frac{\partial z}{\partial z} (f)]$$

$$= -z i\hbar \frac{\partial f}{\partial z} + i\hbar (f) + i\hbar z \frac{\partial f}{\partial z}$$

$$[z, P_z] f = i\hbar f$$

$$\boxed{[z, P_z] = i\hbar}$$

ii) $[x, P_x^n]$

w.k.t $[x, P_x] = i\hbar$

If $n=1 \quad [x, P_x] = i\hbar$

Let $= [x P_x - P_x x] f(x)$

$$= [x (-i\hbar \frac{\partial}{\partial x}) - (-i\hbar \frac{\partial}{\partial x}) x] f(x)$$

$$= -x i\hbar \frac{\partial f(x)}{\partial x} + i\hbar x \frac{\partial f(x)}{\partial x} + i\hbar f(x)$$

$$[x P_x] f(x) = i\hbar f(x)$$

$$[x P_x] = i\hbar$$

if $n=1 \quad [x, P_x^n] \Rightarrow [x, P_x^1] = i\hbar$

by if $n=2 \quad [x, P_x^n] \Rightarrow [x, P_x^2] = 2i\hbar P_x$

if $n=3 \quad [x, P_x^n] \Rightarrow [x, P_x^3] = 3i\hbar P_x^2$

if $n=4$ $[x P_x^n] = [x P_x^4] = 4i\hbar P_x^3$

In general

If $n=n$ $[x P_x^n] = n i\hbar P_x^{n-1}$

$$[x P_x^n] = n i\hbar P_x^{n-1}$$

by $[x P_x^2] \Rightarrow [x, P_x P_x]$

$$= P_x(x P_x) + P_x(x P_x)$$

$$= P_x(i\hbar) + P_x(i\hbar)$$

$$= 2i\hbar P_x$$

IstINTERNAL TEST 20¹² -2023

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Subject/ವಿಷಯ: AMMP

Semester/ಸೆಮಿಸ್ಟರ್: IVth Exam Seat No / ಪರೀಕ್ಷೆ ಸ್ಥಳಾಂಕ:

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1.

Q1) Diffusion or Heat flow equation:

→ the heat flow equation is,

$$\nabla^2 u = \frac{1}{\alpha^2} \frac{\partial u}{\partial t} \quad (1)$$

u = temperature

 ∇^2 = rectangular plate. α = constant

there is flow of the heat is,

$$u = F(x, y, z) T(t) \quad (2)$$

from eqⁿ (1) and (2),

$$\nabla^2 (TF) = \frac{1}{\alpha^2} \frac{\partial (TF)}{\partial t}$$

$$T \nabla^2 F = \frac{1}{\alpha^2} F \frac{\partial T}{\partial t} \quad (3)$$

the eqⁿ (3) divided by (FT)

$$\frac{1}{F} \nabla^2 F = \frac{1}{\alpha^2 T} \frac{\partial T}{\partial t}$$

then, the constant,

$$\frac{1}{F} \nabla^2 F = \frac{1}{\alpha^2 T} \frac{\partial T}{\partial t} = -k^2$$

then separate the equation is

$$\frac{1}{F} \nabla^2 F = -k^2 \Rightarrow \nabla^2 F + k^2 F = 0$$

$$\frac{1}{\alpha^2 T} \frac{\partial T}{\partial t} = -k^2 \Rightarrow \frac{\partial T}{\partial t} + k^2 \alpha^2 T = 0$$

(then, Laplace eqⁿ.)

$$u = e^{-ny} \sin kx$$

$$u = e^{-n} \sin kx$$

$$F = \begin{cases} \sin kx \\ \cos kx \end{cases} \quad T = \begin{cases} e^{-ky} \\ e^{-hy} \end{cases}$$

(at $x=0$ at $y=0$
the $\sin kx = 0$)

$$k = \frac{n\pi}{1}$$

$$k = \frac{n\pi}{x}$$

$$Y = e^{-\frac{\alpha^2 k^2 y^2}{4}} \sin kx$$

$$Y = e^{-\frac{\alpha^2 n^2 \pi^2 y}{4x^2}} \sin kx$$

(2) Wave equation:

the wave equation is,

$$\frac{\partial^2 y}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 y}{\partial t^2} \quad \text{--- (1)}$$

then,

$$(\text{then } y = x) \quad x = y(t) \cdot t(t)$$

$$x = yt$$

$$\frac{\partial^2 yt}{\partial y^2} = \frac{1}{v^2} \frac{\partial^2 (yt)}{\partial t^2} \quad \text{--- (2)}$$

$$yt \frac{\partial^2 y}{\partial y^2} = \frac{1}{v^2} y \frac{\partial^2 t}{\partial t^2} \quad \text{--- (3)}$$

then the eqn (3) divided by yt

$$\frac{1}{y} \frac{\partial^2 y}{\partial y^2} = \frac{1}{v^2 t} \frac{\partial^2 t}{\partial t^2} \quad \text{--- (4)}$$

then constant = k^2

$$\frac{\partial^2 y}{\partial y^2} + k^2 y = 0$$

$$\frac{\partial^2}{\partial t^2} + \sigma k^2 v^2 t = 0$$

then the equation are,

$$X = \begin{Bmatrix} \sin kx \\ \cos kx \end{Bmatrix} \begin{Bmatrix} \sin kvt \\ \cos kvt \end{Bmatrix}$$

$$\text{Or } \bar{X} = \begin{Bmatrix} \sin ky \\ \sin ky \end{Bmatrix} \begin{Bmatrix} \sin k^2 v^2 t \\ \cos k^2 v^2 t \end{Bmatrix}$$

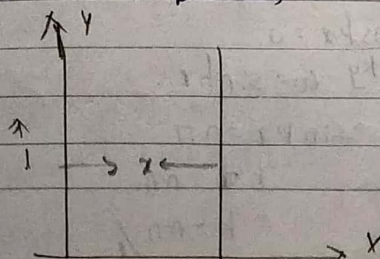
(2) Laplace equation.

Laplace equation: means it is a study state of rectangular plate it also called as. the Laplace "infinite Laplace equation."

then, the infinite

12 We can assume that in a rectangular plate of ~~width~~ so for long separation of the width is extended the y-direction it called "semi finite length"

the rectangular plate,
the length is "l" and width is "x" in the fig shown in the below,



Laplace equation depends on initial temp and final temp. & it determines the plate's width, length and height of the rectangular plate.

the Laplace equation is no flow of the heat.

$$\nabla^2 T = 0$$

(1)

thnd, $\nabla^2 =$ it is the rectangular plate because the side is in rectangular shape.

$$\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} = 0$$

(2)

there, $T = X(x) Y(y)$

X is function of x

Y is function of y .

$$\frac{\partial^2 (XY)}{\partial x^2} + \frac{\partial^2 (XY)}{\partial y^2} = 0$$

or

$$Y \frac{\partial^2 X}{\partial x^2} + X \frac{\partial^2 Y}{\partial y^2} = 0$$

divided by XY

$$\frac{1}{X} \frac{\partial^2 X}{\partial x^2} + \frac{1}{Y} \frac{\partial^2 Y}{\partial y^2} = -k^2$$

$k = \text{constant}$

$$\frac{\partial^2 X}{\partial x^2} + k^2 X = 0 \quad \frac{\partial^2 Y}{\partial y^2} - k^2 Y = 0$$

$$X = \begin{cases} \sin kx \\ \cos kx \end{cases} \quad Y = \begin{cases} e^{+ky} \\ e^{-ky} \end{cases}$$

first at temp = 0 and $y = \infty$

we discard the e^{-ky}

and $t = 0$ $x = 0$

$$\cos kx = 0$$

$$Y = e^{-ky} \sin kx$$

$$x = L, \quad \sin kx = n\pi$$

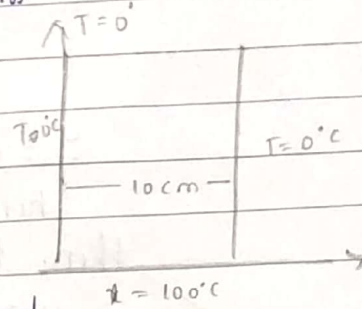
$$kL = n\pi$$

$$k = n\pi/L$$

$$y = e^{-\frac{n\pi y}{1}} \sin kx$$

this is Laplace equation.

- (b) A rectangular metal plate has its 2 sides and
 $T = 0^\circ\text{C}$, $1 = 100^\circ\text{C}$, $2 = 10\text{cm}$
 steady state temp.



this Laplace equation,
 we can assume that in a,
 rectangular plate so for long
 separation of width is extended
 the y-direction it called "semi-finite length"

the Laplace equation. there is no flow of the
 heat.

$$\nabla^2 T = 0$$

(1)

∇^2 is a rectangular plate.

$$\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} = 0$$

(2)

there is, function,

$$T = X(x) Y(y)$$

X is function of x

Y is function of y

$$\frac{\partial^2(XY)}{\partial x^2} + \frac{\partial^2(XY)}{\partial y^2} = 0 \quad (3)$$

$\therefore$ eqn (3) divided by XY

$$\frac{Y}{X} \frac{\partial^2 X}{\partial x^2} + \frac{1}{Y} \frac{\partial^2(Y)}{\partial y^2} = 0$$

$$\frac{1}{X} \frac{\partial^2 X}{\partial x^2} = -k^2 \quad \frac{1}{Y} \frac{\partial^2 Y}{\partial y^2} = -k^2$$

the k = constant

$$\frac{\partial^2 X}{\partial x^2} + k^2 X = 0$$

$$\frac{\partial^2 Y}{\partial y^2} + k^2 Y = 0$$

$$X = \begin{Bmatrix} \sin kx \\ \cos kx \end{Bmatrix}$$

$$Y = \begin{Bmatrix} e^{-ky} \\ e^{ky} \end{Bmatrix}$$

$t=0$ $y=\infty$, first case, they discard
 $t=0$, $x=0$, $\cos kx$ discard

$$y = e^{-ky} \sin kx \quad \text{--- (4)}$$

and $x=10$

$$\sin kx = 0$$

$$k(10) = n\pi$$

$$k = \frac{n\pi}{10}$$

$$y = e^{-\frac{n\pi y}{10}} \sin\left(\frac{n\pi x}{10}\right) \quad \text{--- (5)}$$

here satisfy $t=0$.

then, $t=100$ i.e. $x=10$

$$y = \sum_{n=0}^{\infty} b_n \sin\left(\frac{n\pi x}{10}\right)$$

the sine fourier series,

we find the b_n ,

where $l=10$ $f(x)=100$

$$b_n = \frac{2}{l} \int_0^l f(x) \sin\left(\frac{n\pi x}{l}\right) dx$$

$$b_n = \frac{2}{10} \int_0^{10} 100 \cdot \sin\left(\frac{n\pi x}{10}\right) dx$$

$$\text{the } \begin{cases} \frac{400}{n\pi} & \text{is even} \\ 0 & \text{is odd} \end{cases}$$

$$y = \frac{400}{\pi} \left[e^{-\frac{\pi y}{10}} \sin \frac{n\pi}{10} - e^{-3\pi y/10} \sin \frac{3n\pi}{10} \right] \quad \dots$$

$\frac{\pi y}{10}$ is small, its value variable
 $x=$

$$x=5 \quad y=5$$

$$y = \frac{400}{\pi} \left[e^{\pi/2} \sin(\pi/2) - e^{3\pi/2} \cos\left(\frac{3\pi}{2}\right) \right]$$

$$y = 26.2^\circ$$